

## Superconductivity Studies on Tl-M-Ba-Cu-O ( $M = \text{Ce, Th, Pr, Tb, Pb, and Te}$ ) System<sup>1</sup>

K. A. THOMAS, U. V. VARADARAJU, AND G. V. SUBBA RAO

*Materials Science Research Centre, Indian Institute of Technology,  
Madras 600 036, India*

AND C. V. TOMY AND S. K. MALIK

*Tata Institute of Fundamental Research, Bombay 400 005, India*

Received February 28, 1990

DEDICATED TO J. M. HONIG ON THE OCCASION OF HIS 65TH BIRTHDAY

Superconductivity with a  $T_{c,zero}$  at  $75 \pm 1$  K has been observed in Tl-M-Ba-Cu-O ( $M = \text{Ce, Th}$ ) system for a starting nominal composition of  $\text{Tl}_{2.2}\text{M}_1\text{Ba}_2\text{Cu}_3\text{O}_y$ . X-ray data show that the compounds are multiphasic with the 2201 phase as the predominant one which is responsible for superconductivity. Lower  $T_c$  values are noted for  $M = \text{Pr}$  and  $\text{Tb}$ . Only metallic behavior is observed for pure  $\text{Tl}_2\text{Ba}_2\text{Cu}_1\text{O}_5$  and  $M = \text{Pb}$  or  $(\text{Pb} + \text{Sb})$  or  $\text{Te}$  containing phases down to 15 K. © 1990 Academic Press, Inc.

### Introduction

The discovery of superconductivity in the Tl-Ba-Cu-O system has generated a lot of research activity in thallium-based cuprates. This resulted in the synthesis and characterization of a series of phases having the general formula  $\text{Tl}_m\text{Ba}_2\text{Ca}_{n-1}\text{Cu}_n\text{O}_{m+2(n+1)}$  ( $m = 1$  or  $2$ ;  $n = 1-4$ ) and the property of high temperature superconductivity (HTSC) (1-5). The compound  $\text{Tl}_2\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_{10}$  (2223) possesses the highest known transition temperature ( $T_c$ ) of 125 K. The crystal structures of these phases are layer-type with tetragonal symmetry (4, 5).

<sup>1</sup> Preliminary results have been presented at "Intl. Conf. on Supercond., (ICMAS-89)," Paris, Nov. 23-24, 1989.

The phase with  $m = 2$  and  $n = 1$ , viz.,  $\text{Tl}_2\text{Ba}_2\text{Cu}_1\text{O}_6$  (2201) is a very interesting compound in that it does not contain calcium and the single  $\text{CuO}_2$  plane is sandwiched between the  $\text{BaO}$  layers with the sequence,  $\text{TlO-BaO-CuO}_2\text{-BaO-TlO}$ . In this respect, it resembles the single layer Tl cuprate  $\text{TlBa}_2\text{CuO}_4$  (1201) and  $\text{La}_2\text{CuO}_4$ , which have the layer sequence,  $\text{TlO-BaO-CuO}_2\text{-BaO}$  and  $\text{LaO-CuO}_2\text{-LaO}$ , respectively. The compound 2201 is relatively easy to prepare and has a tetragonal structure with a fairly large  $c$  lattice parameter ( $a = 3.866 \text{ \AA}$ ;  $c = 23.24 \text{ \AA}$ ) due to the fact that it is a bilayer thallium cuprate. An orthorhombically distorted structure ( $a_0 \approx b_0 \approx \sqrt{2}a_1$ ;  $c_0 \approx c_1$ ) can also exist (6, 7). The compound is metallic with low resistivity. The superconducting properties, how-

ever, are at variance and, in the literature,  $T_c$  values ranging from 0 to 80 K have been reported (4–8). Details of preparative conditions, which will affect the oxygen stoichiometry (i.e.,  $y$  in  $Tl_2Ba_2CuO_{6-y}$ ) are assumed to be responsible for the above variations in the reported  $T_c$ . Recent studies by Shimakawa *et al.* (7) indicate that a  $T_{c,zero}$  (zero resistance) of 53 and 50 K respectively are exhibited by both the tetragonal and the orthorhombic polymorphs of 2201 and a decrease in the value of oxygen content ( $\Delta y$ ) of 0.037 can induce superconductivity in a nonsuperconducting phase. The 2201 phase, similar to the other thallium cuprates, is amenable to chemical substitutions at Tl and Ba sites. In addition, several meta or unstable phases can be stabilized by suitable substitution. Thus, Ku and Wu (8) report a  $T_{c,zero}$  of 50 K in  $Tl_2Ba_{2-x}La_xCuO_{6+y}$ ,  $x = 0.2$ . It is now known that many of the HTSCs including the thallium cuprates are near the metal-insulator phase boundary and superconductivity is induced with  $T_c$  increasing smoothly from 0 K near the metal-insulator phase boundary to a maximum value *inside* the metallic region. The phase boundary can be reached either by changes in the oxygen stoichiometry and/or by chemical substitutions.

Recently Wang *et al.* (9) reported a reproducible 90 K superconductivity ( $T_{c,zero} = 80$ –83 K) in thallium cuprates with the nominal starting composition,  $Tl_{2.2}Ba_2Ce_1Cu_3O_y$ . It is interesting to note that compositions with  $Ce > \text{or} < 1.0$  were found to give a lower  $T_c$  (< 20 K). The phase responsible for HTSC has not been clearly identified by Wang *et al.* but they found X-ray lines due to 2201 phase in addition to lines from the impurity phases (CuO and  $BaCeO_3$ ). It is possible that the *proper* 2201 phase gets stabilized by the substitution of  $Tl^{3+}$  by  $Ce^{4+}$  and also the metal-insulator phase boundary is suitably shifted by the optimum concentration of Ce (and thus the oxygen content ( $y$ )) to give rise to HTSC. With this in view,

we have taken up a study of the system  $M$ -Tl-Ba-Cu-O with  $M = Ce, Th, Pr, Tb, Pb,$  and  $Te$ . Our results show that a reproducible  $T_{c,zero}$  of  $75 \pm 1$  K attributable to the 2201 phase can be obtained for  $M_{1.0} = Ce$  and  $Th$ . However,  $Pr$  and  $Tb$  produce a lower  $T_c$  whereas  $Pb, Te,$  and the combination ( $Pb_{0.5} + Sb_{0.5}$ ) as well as the pure 2201 phase do not show a  $T_c$  above 15 K, though all the phases show metallic behavior.

### Experimental

The compounds were prepared using high purity  $Tl_2O_3$  (Cerac, UK: 99.99%),  $BaCO_3$  (Cerac, UK: 99.9%),  $CeO_2$  (IRE, India: 99.99%),  $ThO_2$  (Fluka: 99.9%),  $Sb_2O_3$  (Ventron: 99.99%),  $PbO_2$  (BDH: 99.9%),  $TeO_2$  (Aldrich: 99.9%), and  $CuO$  (Cerac: 99.99%).  $PrO_2$  and  $TbO_2$  were prepared by the solvolytic disproportionation of  $Pr_6O_{11}$  (IRE: 99.99%) and  $Tb_4O_7$  (IRE: 99.99%) in acid solutions and characterized by X-ray diffraction (10). A master composition of  $Ba_2Cu_3O_y$  was prepared using  $BaCO_3$  and  $CuO$  by heating at 900°C in air for a minimum of 48 hr with several intermittent grindings and heatings. Appropriate amounts of  $Tl_2O_3$  and  $MO_2$  (or  $Sb_2O_3$ ) were then added to give a nominal composition,  $Tl_{2.2}MBa_2Cu_3O_y$ . The mixture is then thoroughly ground and pressed into pellets (8 mm diameter, 1–2 mm thick) which are introduced into a preheated tubular furnace (870–900°C) with the oxygen flow. The heating time was 3–5 min. The samples were then pulled out of the hot zone of the furnace and allowed to cool to room temperature by furnace shutoff maintaining the flow of oxygen.

In the  $Th_x$ -Tl-Ba-Cu-O system, compounds with  $x = 0.0, 0.25, 0.50,$  and  $0.75$  have also been prepared. Pure  $Tl_2Ba_2CuO_6$  (2201) phase has also been prepared starting from  $Tl_2O_3, 2BaO_2$  (Cerac: 99.99%), and  $CuO$  (using similar heat-treatment conditions but without recourse to the master composition). Powder X-ray diffraction

(XRD; Philips unit,  $\text{CuK}\alpha$  radiation, Ni filter) has been employed for the phase identification. The lattice parameters were calculated by least-squares (LSQ) fitting of high angle lines using a computer program. Theoretical XRD patterns with the intensities were generated using the Lazypulverix program for comparison with the observed ones. Four probe dc electrical resistivity measurements were made from room temperature down to 68 K using an indigenously fabricated setup with liquid  $\text{N}_2$  or a closed cycle CTI Model 21C He refrigerator down to 15 K. Magnetization measurements were done on select samples using a Quantum Design Model 1822 MPMS SQUID magnetometer.

## Results and Discussion

### a. Structure and Stability

The samples are well-crystalline, black, and are stable under atmospheric conditions. For phases with  $M = \text{Th}$ , we notice slight melting of the samples under the preparative conditions presently employed, especially when  $0.25 < x < 1.0$ . The XRD patterns of all  $\text{Tl}_{2.2}M\text{Ba}_2\text{Cu}_3\text{O}_y$  phases revealed that they are multiphasic. When  $M = \text{Ce}$  and  $\text{Th}$ , the major phase could be indexed as 2201 ( $\text{Tl}_2\text{Ba}_2\text{CuO}_6$ ) and  $\text{CuO}$  was always the impurity phase. In addition  $\text{BaCeO}_3$  and  $\text{ThO}_2$  respectively were also observed as the other impurity phases (Fig. 1). Wang *et al.* (9) also noted  $\text{BaCeO}_3$  and  $\text{CuO}$  impurity phases in their studies. In the  $\text{Tl}_{2.2}\text{Th}_x\text{Ba}_2\text{Cu}_3\text{O}_y$  series of compounds we noted a decrease in the intensity of the  $\text{ThO}_2$  lines as  $x$  decreases from 1.0 to 0.25. On the other hand, when  $M = \text{Pr}$  and  $\text{Tb}$ , the major phase lines correspond to 1212 type and not the 2201 type structure. When  $M = \text{Pb}$ , unequivocal indexing of the XRD patterns was not possible.

The composition  $\text{Tl}_{2.2}\text{Ba}_2\text{Cu}_3\text{O}_y$  also gave  $\text{Tl}_2\text{Ba}_2\text{Cu}_3\text{O}_6$  (2201) as the major phase with  $\text{CuO}$  as impurity whereas  $\text{Tl}_2\text{Ba}_2\text{CuO}_6$  pre-

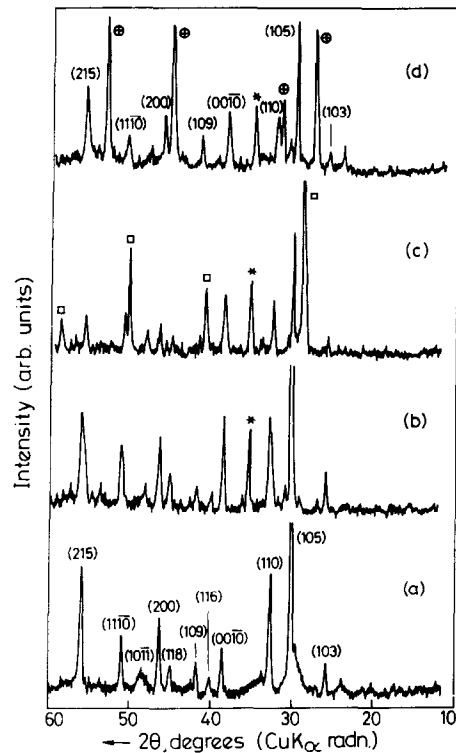


FIG. 1. Powder X-ray diffraction patterns of the phases with nominal compositions in the Tl-M-Ba-Cu-O system: (a)  $\text{Tl}_2\text{Ba}_2\text{Cu}_3\text{O}_y$ , (b)  $\text{Tl}_{2.2}\text{Ba}_2\text{Cu}_3\text{O}_y$ , (c)  $\text{Tl}_{2.2}\text{CeBa}_2\text{Cu}_3\text{O}_y$ , (d)  $\text{Tl}_{2.2}\text{ThBa}_2\text{Cu}_3\text{O}_y$ . The patterns can be indexed as the  $\text{Tl}_2\text{Ba}_2\text{CuO}_6$  (2201) phases with minor impurity phases. (Impurity lines due to \*,  $\text{CuO}$ ; □,  $\text{BaCeO}_3$ ; ⊕,  $\text{ThO}_2$ .)

pared from  $\text{Tl}_2\text{O}_3$ ,  $\text{BaO}_2$ , and  $\text{CuO}$  (nominal 2201) formed a single phase material (Fig. 1). The tetragonal lattice parameters presently obtained by LSQ fitting of  $(hkl)$  lines are  $a = 3.87 \text{ \AA}$  and  $c = 23.21 \text{ \AA}$  and are in good agreement with those reported in the literature (5).

### b. Electrical and Superconducting Properties

All the phases in the system Tl-M-Ba-Cu-O including the pure 2201 phase were found to be metallic with low room temperature resistivities ( $\rho_{300 \text{ K}} =$

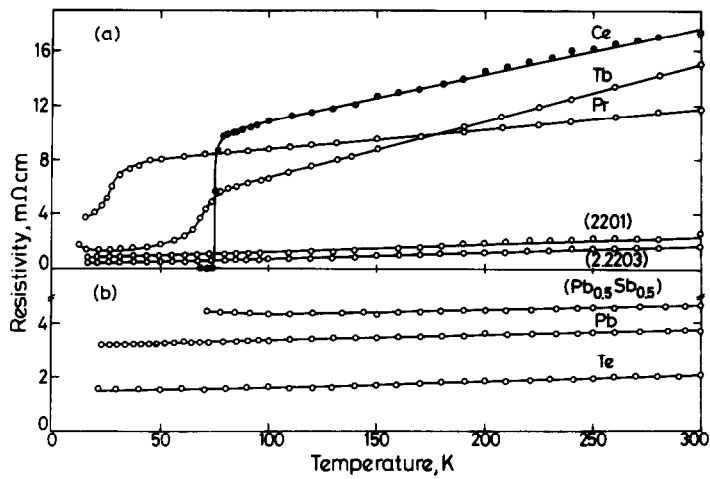


FIG. 2. Resistivity vs temperature plots for the phases with nominal composition,  $Tl_{2.2}M_1Ba_2Cu_3O_y$ ; (a)  $M = Ce, Pr, \text{ and } Tb$ ; (b)  $M = Pb, Te \text{ and } (Sb_{0.5}Pb_{0.5})$ . Data on nominal  $Tl_{2.2}Ba_2Cu_3O_6$  also included (a). As can be seen, the phase with  $M = Ce$  exhibits HTSC.

1–35  $m\Omega \text{ cm}$ ). Superconductivity is encountered for  $M = Ce, Th, Pr, \text{ and } Tb$ . The Ce-containing phase showed a  $T_{c,zero}$  value of 74 K (Fig. 2a) which is less than that reported (80–83 K) by Wang *et al.* (9). The Th-containing phase, presently synthesized for the first time, showed a  $T_{c,zero}$  of 75 K (Fig. 3).

The  $T_{c,onset}$  values are  $\sim 80$  and  $\sim 90$  K respectively for  $M = Ce$  and Th phases.

Superconductivity is also shown for  $x = 0.25, 0.50, \text{ and } 0.75$  in the  $Tl_{2.2}Th_xBa_2Cu_3O_y$  system. Even though the  $T_{c,onset}$  (taken as the temperature at which significant departure from the linear  $\rho$ - $T$  behavior is noted)

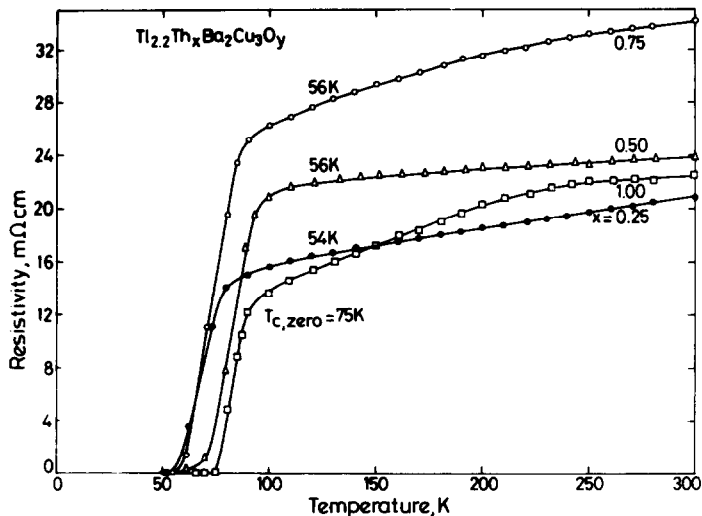


FIG. 3. Resistivity vs temperature plots for  $Tl_{2.2}Th_xBa_2Cu_3O_y$ , showing superconductivity for  $x \geq 0.25$ . Values of  $x$  and  $T_{c,zero}$  are indicated.

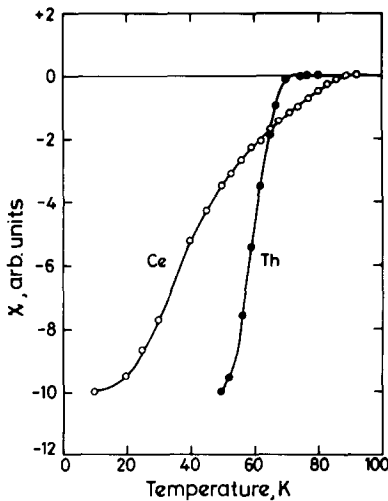


FIG. 4. Magnetic susceptibility (arbitrary units) vs temperature plots for the phases with nominal compositions  $Tl_{2.2}MBa_2Cu_3O_y$ ;  $M = Ce$  and  $Th$  showing superconducting transitions.

values are in the 75–100 K region, the  $T_{c,zero}$  values are lower. It is interesting, however, that  $T_{c,zero}$  values are clustered in a narrow range (54–56 K) for  $x < 0.75$  (Fig. 3). This behavior is similar to that noted by Wang *et al.* (9) for the  $Tl_{2.2}Ce_xBa_2Cu_3O_y$  system ( $x \leq 1.0$ ) and conclusively establish that a nominal concentration of  $M_{1.0}$  ( $M = Ce$  or  $Th$ ) is necessary to stabilize the 80 K superconductor.

Magnetic susceptibility studies done on  $Tl_{2.2}MBa_2Cu_3O_y$  ( $M = Ce$  and  $Th$ ) phases show that the superconductivity is bulk in nature (Fig. 4). The  $T_c$  value of 75 K ( $M = Th$ ) determined from the temperature of diamagnetic onset agrees excellently with the  $T_{c,zero}$  obtained by  $\rho$ - $T$  measurements. For  $M = Ce$ , the  $T_c$  value of 85 K from  $\chi$ - $T$  data is larger than the  $T_{c,zero}$  (74 K) from  $\rho$ - $T$  data but compares well with that reported by Wang *et al.* (9). We also note that the transition with  $M = Ce$  is broader compared to that with  $M = Th$ .

The 2201 phases obtained from nominal compositions  $Tl_{2.2}Ba_2Cu_3O_y$  and  $Tl_2Ba_2CuO_6$  (see Experimental) were found to be

only metallic and no superconductivity was encountered down to 15 K, the lowest temperature reached with our set up. Our findings are in agreement with those of several workers in the literature (4, 7, 8, 11). For the phases with  $M = Pr$  and  $Tb$  a clear drop in resistivity was noted at 35 and 75 K respectively indicating the onset of superconductivity (Fig. 2a). However, zero resistance has not been observed down to 15 K. For  $M = Pb$ ,  $(Pb_{0.5}Sb_{0.5})$  and  $Te$  in the  $Tl_{2.2}MBa_2Cu_3O_y$  system, an almost temperature-independent resistivity behavior is noticed with no superconducting transition down to 15 K (Fig. 2b).

The room temperature resistivities of the phases which show superconductivity were found to be three–five times larger than non-superconducting but metallic phases, including the pure 2201 phase. This indicates the approach toward the metal-to-insulator phase boundary in the  $M$ -substituted compounds. However, the role of  $Ce$  and  $Th$  ions in the stabilization of HTSC phase is not clear. Also the importance of a nominal composition of  $M_{1.0}$  is intriguing even though XRD clearly reveals that only part of  $Ce$  or  $Th$  is entering into the lattice, possibly at the Tl site. It is not clear whether doping by  $M$  ion is directly producing a decrease in charge carrier concentration (approaching insulator or semiconductor behavior) or indirectly affecting the oxygen stoichiometry in the  $Tl_2Ba_2Cu_1O_{6+y}$  phase. Studies by Rao *et al.* (12) and Parise *et al.* (13) clearly showed that  $Pb$  substitutes at the Tl site as  $Pb$  ion in  $(Tl, Pb)$ - $Sr$ - $Ca$ - $Cu$ - $O$  systems. Wang *et al.* (9) mentioned that their preliminary Seebeck data indicate an  $n$ -type behavior for the  $Tl_{2.2}Ce_1Ba_2Cu_3O_y$  phase indicating that these are electron superconductors.

As mentioned earlier in the case of  $Pr$  and  $Tb$ -containing phases the major phase appears to be 1212 from XRD studies. In such a case, the phase should exhibit semiconductor behavior (14). However, it is possible that a second phase exists similar to

that reported by Bourgault *et al.* (15) wherein partial substitution at the Tl and Sr sites simultaneously has been suggested with the composition  $Tl_{1-x}Pr_xSr_{2-y}Pr_yCuO_z$ . It is further suggested that Pr enters the lattice as  $Pr^{4+}$  at the Tl site and as  $Pr^{3+}$  at the Sr site. In the present case similar substitution at the Tl and Ba sites (instead of Sr) may be taking place. The observed metallic behavior and onset of superconductivity could be due to such a phase (the  $T_c$  value reported for the Sr compound is 30 K from  $\chi$ - $T$  measurements (15)).

Detailed Seebeck and Hall data will indeed throw light on the nature of the Ce- and Th-containing phases and HTSC. These studies are underway.

### Conclusions

Our results show that superconductivity with  $T_{c,zero}$  of  $75 \pm 1$  K can be achieved in a starting nominal composition of the system  $Tl_{2.2}M_1Ba_2Cu_3O_y$ , for  $M = Ce$  and  $Th$ , but not when  $M = Pb$  or  $Te$  under the preparative conditions presently employed. Nominal  $M_{1.0}$  is necessary to stabilize the 75 K phase and the  $T_{c,zero}$  decreases to 54 K when  $M = Th$  and  $x \leq 0.75$ . Phases where  $M = Pr$  and  $Tb$  also exhibit a drop in resistivity below 70 K, but zero resistance is not achieved.

### Acknowledgments

Thanks are due to PMB on superconductivity, D.S.T. Government of India, New Delhi and Indo-EEC Project for the award of research grants.

### References

1. Z. Z. SHENG AND A. M. HERMAN, *Nature (London)* **332**, 55 (1988); **332**, 138 (1988).
2. C. N. R. RAO (Ed.), "Chemical and Structural Aspects of High Temperature Superconductors," World Scientific, Singapore (1988).
3. C. N. R. RAO AND B. RAVEAU, *Acc. Chem. Res.* **22**, 106 (1989); T. V. Ramakrishnan and C. N. R. Rao, *J. Phys. Chem.* **93**, 4414 (1989).
4. S. S. P. PARKIN, V. Y. LEE, A. I. NAZZAL, R. SAVOY, T. C. HUANG, G. GORMAN, AND R. BEYERS, *Phys. Rev. B* **38**, 6531 (1988).
5. C. C. TORARDI, M. A. SUBRAMANIAN, J. C. CALABRESE, J. GOPALAKRISHNAN, E. M. MCCARRON, K. J. MORRISSEY, T. R. ASKEW, R. B. FLIPPEN, U. CHOWDHRY, AND A. W. SLEIGHT, *Phys. Rev. B* **38**, 225 (1988); D. E. COX, C. C. Torardi, M. A. Subramanian, J. Gopalakrishnan, and A. W. Sleight, *Phys. Rev. B* **38**, 6624 (1988).
6. J. B. PARISE, J. GOPALAKRISHNAN, M. A. SUBRAMANIAN, AND A. W. SLEIGHT, *J. Solid State Chem.* **76**, 432 (1988).
7. Y. SHIMAKAWA, Y. KUBO, T. MANAKO, T. SATOH, S. IJIMA, T. ICHIHASHI, AND H. IGARASHI, *Physica C* **157**, 279 (1989).
8. H. C. KU AND P. T. WU (Eds.), "Progress in High Tc Superconductivity," Proc. of Taiwan Intl. Symp. on Supercond., Vol. 19, World Scientific, Singapore (1989).
9. J. H. WANG, Z. Z. SHENG, C. DONG, X. FEI, L. SHENG, A. M. HERMANN, AND Z. X. ZHAO, *Physica C*, **158**, 507 (1989).
10. R. L. N. SASTRY, P. N. MEHROTRA, AND C. N. R. RAO, *J. Inorg. Nucl. Chem.* **28**, 2116 (1966).
11. M. KIKUCHI, T. KAKITANI, T. SUZUKI, S. NAKAJIMA, K. KIRAGA, N. KOBAYASHI, H. IWASAKI, Y. SYONO, AND Y. MUTO, *Japan J. Appl. Phys.* **28**, L382 (1989).
12. C. N. R. RAO, A. K. GANGULI, AND R. VJAYARAGHAVAN, *Phys. Rev. B* **40**, 2565 (1989).
13. J. B. PARISE, P. L. GAI, M. A. SUBRAMANIAN, J. GOPALAKRISHNAN, AND A. W. SLEIGHT, *Physica C* **159**, 245 (1989).
14. C. MARTIN, D. BOURGALT, C. MICHEL, M. HERVIEU, AND B. RAVEAU, *Mod. Phys. Lett. B* **3**, 93 (1989).
15. D. BOURGALT, C. MARTIN, C. MICHEL, M. HERVIEU, J. PROVOST, AND B. RAVEAU, *J. Solid State Chem.* **78**, 326 (1989).